



Dynkin Diagrams of Simple Lie Algebras



$B_2(3)$	$C_2(4)$	$D_2(2)$	${}^2D_2(2)$
$B_2(4)$	$C_2(5)$	$D_2(3)$	${}^2D_2(3)$

A_{1-6}
1
1

A_{1-6}	A_{1-6}
A_1	$A_1(7)$
A_2	A_2
A_{1-6}	A_{1-6}
A_4	$A_1(8)$
A_4	A_4

A_7	$A_1(11)$	$E_6(2)$	$E_7(2)$	$E_8(2)$	$E_8(2)$	$G_2(3)$	${}^2D_4(2)$	${}^2F_4(2)$	${}^2B_2(2)$	${}^2F_4(2)'$	${}^2G_2(2)$	$B_2(2)$	$C_2(3)$	$D_2(2)$	${}^2D_2(2)$
A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7	A_7

A_{10}	$A_2(13)$	$F_4(3)$	$E_5(3)$	$F_4(3)$	$E_6(3)$	$G_2(4)$	${}^2D_4(3)$	${}^2F_4(3)$	${}^2B_2(3)$	${}^2F_4(3)$	${}^2G_2(3)$	$B_2(3)$
A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}	A_{10}

A_{14}	$A_1(17)$	$E_6(4)$	$E_7(4)$	$E_8(4)$	$E_8(4)$	$G_2(5)$	${}^2D_4(4)$	${}^2F_4(4)$	${}^2B_2(4)$	${}^2F_4(4)$	${}^2G_2(4)$	$B_2(4)$
A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}	A_{14}

A_{18}	$A_1(19)$	$F_4(5)$	$E_7(5)$	$F_4(5)$	$E_8(5)$	$G_2(6)$	${}^2D_4(5)$	${}^2F_4(5)$	${}^2B_2(5)$	${}^2F_4(5)$	${}^2G_2(5)$	$B_2(5)$
A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}	A_{18}

A_{22}	$A_1(21)$	$F_4(6)$	$E_8(6)$	$F_4(6)$	$E_9(6)$	$G_2(7)$	${}^2D_4(6)$	${}^2F_4(6)$	${}^2B_2(6)$	${}^2F_4(6)$	${}^2G_2(6)$	$B_2(6)$
A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}	A_{22}

- Chevalley Groups
- Classical Chevalley Groups
- Classical Groups
- Classical Orthogonal Groups
- Matrix Groups
- Sporadic Groups
- "Near Groups and This Group"
- Simple Groups
- Sporadic Groups
- Finite Groups

Minimal Symbol	Group
M_{11}	M_{11}
M_{12}	M_{12}
M_{13}	M_{13}
M_{14}	M_{14}
J_1	J_1
J_2	J_2
J_3	J_3
J_4	J_4
J_5	J_5
J_6	J_6
J_7	J_7
J_8	J_8
J_9	J_9
J_{10}	J_{10}
J_{11}	J_{11}
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J_{99}	J_{99}
J_{100}	J_{100}

Neuigkeiten über Lie-Algebren

Kolloquium Tübingen

17. Juli 2017



Meinolf Geck

The greatest mathematical paper of all time

A. J. Coleman, *The Mathematical Intelligencer* **11** (1989), 29–38

Euclid's *Elements*, Newton's *Principia*, . . . and in the past 200 years:

Wilhelm Killing (1847–1923), *Die Zusammensetzung der stetigen, endlichen Transformationsgruppen II*, *Mathematische Annalen* **33** (1888), 1–48.

MathSciNet Review MR1007036 by Jean Dieudonné:

“Many mathematicians will disagree, [...]. One may however observe that:

- (1) nobody had tackled the problem before Killing;
- (2) he solved it by methods he invented;
- (3) nobody collaborated with him for that solution;
- (4) the result became a most important milestone in modern mathematics.

I think it is not so easy to find papers exhibiting all those features”

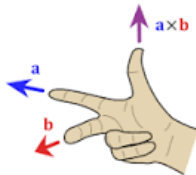
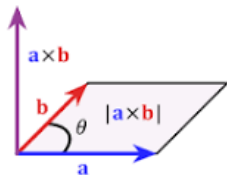
Lie algebra = infinitesimal version of a transformation group

A vector space L equipped with a product $x \cdot y$ such that

- $x \cdot x = 0$,
- $x \cdot (y \cdot z) + y \cdot (z \cdot x) + z \cdot (x \cdot y) = 0$ (Jacobi identity).

Usually write $x \cdot y$ as $[x, y]$ (Lie bracket).

Example. Vector product in $L = \mathbb{R}^3$:



Example. Matrices: $L = M_n(\mathbb{C})$ with bracket $[A, B] = A \cdot B - B \cdot A$; denote $\mathfrak{gl}_n(\mathbb{C})$.
Subalgebra $\mathfrak{sl}_n(\mathbb{C}) = \{A \in M_n(\mathbb{C}) \mid \text{Trace}(A) = 0\}$.

An infinite-dimensional example

Start with $R = \mathbb{C}[X, X^{-1}]$ (Laurent polynomials).

A linear map $D: R \rightarrow R$ is called a derivation if

$$D(f \cdot g) = f \cdot D(g) + D(f) \cdot g \quad \text{for all } f, g \in R.$$

Example: $D =$ usual (formal) derivative with respect to X .

$$D(X^n) = nX^{n-1}, \quad D(X^{-1}) = -X^{-2}, \quad \text{etc.}$$

Let $L =$ vector space of all derivations of R .

Exercise: If $D, D' \in L$, then $[D, D'] = D \circ D' - D' \circ D \in L$. So L is a Lie algebra.

For $m \in \mathbb{Z}$, we set $L_m(f) = -X^{m+1} \cdot D(f)$ for $f \in R$. Then:

- $\{L_m \mid m \in \mathbb{Z}\}$ is a vector space basis of L .
- $[L_m, L_n] = (m - n)L_{m+n}$ for all $m, n \in \mathbb{Z}$.

L is called “Witt algebra” \rightsquigarrow important in mathematical physics.

A subspace $U \subseteq L$ is called an ideal if $[u, x] \in U$ and $[x, u] \in U$ for all $x \in L, u \in U$. In this case, $\bar{L} = L/U$ also is a Lie algebra. So L “built up” from U and \bar{L} .

Definition. L is called simple if $L \neq \{0\}$, the bracket is not identically zero and there is no proper ideal in L .

Cartan–Killing (~1890): The finite-dimensional simple Lie algebras over \mathbb{C} are classified by “Dynkin diagrams”.



Infinite families: Lie algebras of matrices

$$A_n \leftrightarrow \mathfrak{sl}_{n+1}(\mathbb{C}), \quad B_n \leftrightarrow \mathfrak{so}_{2n+1}(\mathbb{C}), \quad C_n \leftrightarrow \mathfrak{sp}_{2n}(\mathbb{C}), \quad D_n \leftrightarrow \mathfrak{so}_{2n}(\mathbb{C}).$$

Exceptional algebras:

$$\dim \mathfrak{g}_2 = 14, \quad \dim \mathfrak{f}_4 = 52, \quad \dim \mathfrak{e}_6 = 78, \quad \dim \mathfrak{e}_7 = 133, \quad \dim \mathfrak{e}_8 = 248.$$

S. GARIBALDI, E_8 , the most exceptional group.

Bull. AMS **53** (2016), 643–671.

“The Lie algebra \mathfrak{e}_8 or Lie group E_8 was first sighted by a human being sometime in summer or early fall of 1887, by Wilhelm Killing as part of his program to classify the semisimple finite-dimensional Lie algebras over the complex numbers. [...]

Since then, it has been a source of fascination for mathematicians and others in its role as the largest of the exceptional Lie algebras. [...]

If we know some statement for all groups except E_8 , then we do not really know it.”

How does a Dynkin diagram determine a simple Lie algebra ?

If diagram has nodes $1, \dots, n$, then form “Cartan matrix” $A = (a_{ij})_{1 \leq i, j \leq n}$ where:

- $a_{ii} = 2$ for all i ; $a_{ij} = 0$ if $i \neq j$ are not joined by an edge;
- if $i \neq j$ are joined by a single edge, set $a_{ij} = a_{ji} = -1$;
- if $i \neq j$ are joined by $m \geq 2$ edges (arrow towards j), set $a_{ij} = -1$, $a_{ji} = -m$.

$$B_3 \quad \begin{array}{c} 1 \quad 2 \quad 3 \\ \bullet \leftarrow \bullet \text{---} \bullet \end{array} \quad \rightsquigarrow \quad A = \begin{pmatrix} 2 & -2 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 2 \end{pmatrix}$$

Chevalley generators (“épinglage”) of corresponding Lie algebra L :

- $L = \langle e_i, f_i, h_i \mid 1 \leq i \leq n \rangle_{\text{Lie}}$,
- $[e_i, f_i] = h_i$, $[e_i, f_j] = 0$ (if $i \neq j$), $[h_i, e_j] = a_{ij}e_j$, $[h_i, f_j] = -a_{ji}f_j$.
- $H := \langle h_1, \dots, h_n \rangle_{\mathbb{C}}$ “Cartan” subalgebra of L .

This determines L up to isomorphism.

Vector space basis of L

Consider \mathbb{R}^n with standard basis $\{\alpha_1, \dots, \alpha_n\}$. Define

$$s_j: \mathbb{R}^n \rightarrow \mathbb{R}^n, \quad \alpha_j \mapsto \alpha_j - a_{ij}\alpha_i.$$

Then $s_j^2 = \text{id}_n$ and $W = \langle s_1, \dots, s_n \rangle \subseteq \text{GL}_n(\mathbb{R})$ is a finite reflection group.

$\Phi := \{w(\alpha_i) \mid w \in W, 1 \leq i \leq n\}$ is the abstract “root system” associated with A .
Corresponding Lie algebra L has basis $\{h_1, \dots, h_n\} \cup \{e_\alpha \mid \alpha \in \Phi\}$.

“Root strings” (Killing): Let $\alpha, \beta \in \Phi$, $\alpha \neq \pm\beta$.

Let $p = p_{\alpha,\beta} \geq 0$ and $q = q_{\alpha,\beta} \geq 0$ be maximal such that

$$\beta - q\alpha, \quad \dots, \quad \beta - \alpha, \quad \beta, \quad \beta + \alpha, \quad \dots, \quad \beta + p\alpha \in \Phi.$$

C. CHEVALLEY (1955): e_α can be chosen such that $e_{\alpha_i} = \pm e_i$, $e_{-\alpha_i} = \pm f_i$ and

$$[e_\alpha, e_\beta] = \pm(q_{\alpha,\beta} + 1)e_{\alpha+\beta} \quad \text{whenever } \alpha, \beta, \alpha + \beta \in \Phi$$

(and with this normalization, the e_α are unique up to sign).

An aside: root systems and lattices

Let $\Gamma :=$ all \mathbb{Z} -linear combinations of roots $\alpha \in \Phi$. Then Γ is a “lattice” in \mathbb{R}^n .

Conversely, let $\Gamma \subseteq \mathbb{R}^n$ be a lattice.

- Γ is **integral** if $(x, y) \in \mathbb{Z}$ for $x, y \in \Gamma$; then Γ is **even** if $(x, x) \in 2\mathbb{Z}$ for $x \in \Gamma$.
- **Dual lattice** $\Gamma^* := \{x \in \mathbb{R}^n \mid (x, y) \in \mathbb{Z} \text{ for all } y \in \Gamma\}$; if Γ integral, then $\Gamma \subseteq \Gamma^*$.
- Γ is **unimodular** if Γ is integral and $\Gamma = \Gamma^*$.

Up to isomorphism, there is only one even unimodular lattice in \mathbb{R}^8 .

(History goes back to around 1870, before Lie and Killing.) This is given by

$$\Gamma_8 := \{(x_i) \in \mathbb{R}^8 \mid 2x_i \in \mathbb{Z}, x_i - x_j \in \mathbb{Z}, \sum_i x_i \in 2\mathbb{Z}\}.$$

The set $R := \{x \in \Gamma_8 \mid (x, x) = 2\}$ contains precisely 240 vectors; these form the “root system” associated with the Dynkin diagram E_8 .

(A further famous example in dimension 24: the Leech lattice \rightsquigarrow Conway’s finite simple groups.

See, e.g., J-P. Serre, A Course in Arithmetic, or W. Ebeling, Lattices and Codes, Springer-Verlag.)

C. CHEVALLEY (~1955/1960): Algebraic analogues of Lie groups.

Simple Lie algebra L + field k \rightsquigarrow group $G_L(k)$.

(Mimic exponential “ $\exp(e_\alpha)$ ” over arbitrary field k ; then $G_L(k) = \langle \exp(te_\alpha) \mid \alpha \in \Phi, t \in k \rangle$.)

- k algebraically closed: $G_L(k)$ simple algebraic group over k ,
 $SL_n(k)$, $Sp_{2n}(k)$, $SO_n(k)$, $E_8(k)$, ...
- k finite \rightsquigarrow new families (at the time) of finite simple groups,
e.g., $|E_8(\mathbb{F}_q)| = q^{248} + \text{lower powers of } q$ (where q is a prime power).

Textbook references:

- BOURBAKI, Groupes et algèbres de Lie, 1968/1975;
- STEINBERG, Lectures on Chevalley groups, 1967/68 (now available from AMS);
- HUMPHREYS, Introduction to Lie algebras and representation theory, 1972;
- CARTER, Simple groups of Lie type, 1972;
- ERDMANN–WILDON, Introduction to Lie algebras, 2006.

So what are the news about Lie algebras ?



GEORGE LUSZTIG,
arXiv:1309.1382 (3 pages):

“The Lie group of type E_8
can be obtained from the
graph E_8 by a method of
Chevalley (1955),

simplified using theory
of canonical basis (1990).”

↪ 3 more papers:
one by Lusztig, two by myself.

Why look for a simplification ?

- Construction of L is an issue, especially for exceptional types.
- Chevalley's construction of $G_L(k)$ relies on choice of e_α 's.

J. TITS 1966: Start with vector space M of correct dimension, fix basis $\{h_1, \dots, h_n\} \cup \{e_\alpha \mid \alpha \in \Phi\}$ and try to define Lie bracket for basis elements. Relies on systematic study of sign choices in Chevalley's normalisation of e_α .

J-P. SERRE 1966: Start with free Lie algebra generators $\{e_i, f_i \mid 1 \leq i \leq n\}$ and add defining relations (nowadays called "Serre relations") to obtain a presentation. Very elegant, does not resolve issue of making choices for the e_α 's.

C.-M. RINGEL 1990: Fix orientation of Dynkin diagram and use the representation theory of quivers and Hall polynomials. Then e_α 's correspond to well-defined indecomposable objects in certain category of modules. Pre-cursor of "canonical bases", still relies on 2^{n-1} choices of orientations.

The simplified construction

Recall “root strings”: Let $\alpha, \beta \in \Phi$, $\alpha \neq \pm\beta$. Then

$$p_{\alpha, \beta} := \max\{i \geq 0 \mid \beta + i\alpha \in \Phi\} \quad \text{and} \quad q_{\alpha, \beta} := \max\{j \geq 0 \mid \beta - j\alpha \in \Phi\}.$$

Let M be a vector space over \mathbb{C} of the “correct” dimension, with a given basis $\{u_1, \dots, u_n\} \cup \{v_\alpha \mid \alpha \in \Phi\}$. Define linear maps $e_i, f_i: M \rightarrow M$ by

$$e_i(u_j) := |a_{ji}|v_{\alpha_j}, \quad e_i(v_\alpha) := \begin{cases} (q_{\alpha_j, \alpha} + 1)v_{\alpha + \alpha_j} & \text{if } \alpha + \alpha_j \in \Phi, \\ u_j & \text{if } \alpha = -\alpha_j, \\ 0 & \text{otherwise;} \end{cases}$$
$$f_i(u_j) := |a_{ji}|v_{-\alpha_j}, \quad f_i(v_\alpha) := \begin{cases} (p_{\alpha_j, \alpha} + 1)v_{\alpha - \alpha_j} & \text{if } \alpha - \alpha_j \in \Phi, \\ u_j & \text{if } \alpha = \alpha_j, \\ 0 & \text{otherwise.} \end{cases}$$

(Note: Matrices of e_i, f_i have entries in $\mathbb{Z}_{\geq 0}$ — in fact, in $\{0, 1, 2, 3\}$.)

Theorem. LUSZTIG (1990) + G. (Proc. Amer. Soc., 2017)

Consider the Lie algebra $\mathfrak{gl}(M)$ and let $L := \langle e_i, f_i \mid 1 \leq i \leq n \rangle_{\text{Lie}} \subseteq \mathfrak{gl}(M)$.

Then L is a simple Lie algebra corresponding to the given Dynkin diagram.

Example: Type G_2 with Cartan matrix $A = \begin{pmatrix} 2 & -1 \\ -3 & 2 \end{pmatrix}$.

- \mathbb{R}^2 with standard basis $\{\alpha_1, \alpha_2\}$. Matrices $s_1, s_2 \in \text{GL}_2(\mathbb{R})$ given by

$$s_1 = \begin{pmatrix} -1 & 1 \\ 0 & 1 \end{pmatrix}, \quad s_2 = \begin{pmatrix} 1 & 0 \\ 3 & -1 \end{pmatrix}.$$

- Applying s_1, s_2 repeatedly to $\alpha_1 = (1, 0)$, $\alpha_2 = (0, 1)$, we obtain

$$\Phi = \{\pm(1, 0), \pm(0, 1), \pm(1, 1), \pm(1, 2), \pm(1, 3), \pm(2, 3)\}.$$

- Simple Lie algebra $L = \langle e_1, e_2, f_1, f_2 \rangle_{\text{Lie}} \subseteq \mathfrak{gl}_{14}(\mathbb{C})$. For example:

$$e_1 = \begin{pmatrix} \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & 2 & 3 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 \end{pmatrix}$$

e_1, e_2 upper triangular

f_1, f_2 lower triangular

(A dot stands for 0)

Resolving sign issue in Chevalley's basis

We can transport the basis $\{u_1, \dots, u_n\} \cup \{v_\alpha \mid \alpha \in \Phi\}$ from M to L and obtain:

Corollary.

- There is a basis $\{\tilde{h}_1, \dots, \tilde{h}_n\} \cup \{e_\alpha \mid \alpha \in \Phi\}$ of L such that:

$$[e_i, \tilde{h}_j] = |a_{ji}| e_{\alpha_i}, \quad [e_i, e_\alpha] = \begin{cases} (q_{\alpha_i, \alpha} + 1) e_{\alpha + \alpha_i} & \text{if } \alpha + \alpha_i \in \Phi, \\ \tilde{h}_j & \text{if } \alpha = -\alpha_i, \\ 0 & \text{otherwise;} \end{cases}$$
$$[f_i, \tilde{h}_j] = |a_{ji}| e_{-\alpha_i}, \quad [f_i, e_\alpha] = \begin{cases} (p_{\alpha_i, \alpha} + 1) e_{\alpha - \alpha_i} & \text{if } \alpha - \alpha_i \in \Phi, \\ \tilde{h}_j & \text{if } \alpha = \alpha_i, \\ 0 & \text{otherwise.} \end{cases}$$

- This basis of L is unique up to a global constant.
- In particular, the structure constants in the equations $[e_\alpha, e_\beta] = N_{\alpha\beta} e_{\alpha+\beta}$ (for $\alpha, \beta, \alpha + \beta \in \Phi$) are uniquely determined up to a global constant.

Where do the formulae for the action of e_i, f_i come from?

- \approx 1985: Introduction of **quantised enveloping algebras** of Lie algebras (Drinfeld, Jimbo), as Hopf algebra deformations depending on a parameter v . Origins in mathematical physics, quantum integrable systems, . . .
- Early 1990's: Discovery of "**canonical bases**" (Lusztig) and "**crystal bases**" (Kashiwara). Specialisation $v \mapsto 1$ gives rise to canonical bases in irreducible representations of ordinary simple Lie algebras.

LUSZTIG (1990 + J. Comb. Algebra 2017). Consider $L \hookrightarrow \mathfrak{gl}(M)$, as above. This is an irreducible representation and $\{u_1, \dots, u_n\} \cup \{v_\alpha \mid \alpha \in \Phi\}$ is the "**canonical basis**" of M in the above sense. (\rightsquigarrow Explains why formulae are "natural".)

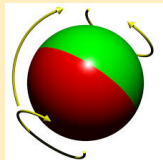
Explicit/simple construction of L and, hence, also of $G_L(k)$, useful for:

- algorithmic problems: nilpotent orbits, matrix group recognition project, . . .;
- teaching courses on Lie algebras ("existence theorem").

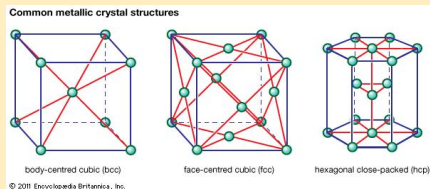
My motivation for studying Lie algebras and Chevalley groups

Group theory = study of symmetries

Continuous \rightsquigarrow Lie groups



Discrete \rightsquigarrow finite groups

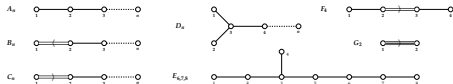


- “Atoms” of symmetry: finite **simple** groups.
- A highlight of 20th century mathematics: **Classification**.
(first announced 1981 • completed 2004: Aschbacher, Smith • 12000 pages of proof)
- 2nd generation proof: D. Gorenstein, R. Lyons, R. Solomon, Math. Survey and Monographs, Amer. Math. Soc., 1994 –?? (currently 6 volumes).

The Periodic Table Of Finite Simple Groups

B, C, Z_1
1
1

Dynkin Diagrams of Simple Lie Algebras



$A_1(4), A_1(5)$	$A_2(2)$
A_5	$A_1(7)$
60	168
$A_1(9), B_2(2)'$	${}^2G_2(3)'$
A_6	$A_1(8)$
360	504

${}^2A_3(4)$	$C_3(3)$	$D_4(2)$	${}^2D_4(2^2)$	$G_2(2)'$
$B_2(3)$	$C_3(3)$	$D_4(2)$	${}^2D_4(2^2)$	${}^2A_2(9)$
25920	436832400	174382400	197406720	6408
$B_2(4)$	$C_3(5)$	$D_4(3)$	${}^2D_4(3^2)$	${}^2A_2(16)$
979200	2385018000000	493217910400	101311968419520	62400

A_7	$A_1(11)$	$E_6(2)$	$E_7(2)$	$E_8(2)$	$F_4(2)$	$G_2(3)$	${}^3D_4(2^3)$	${}^2F_4(2^2)$	${}^2B_2(2^2)$	${}^2F_4(2)'$	${}^2G_2(3^3)$	$B_3(2)$	$C_4(3)$	$D_5(2)$	${}^2D_5(2^2)$	${}^2A_2(25)$	
2520	660	214861375322	7984080000000000	4032000000000000000	3331326	4245696	231341312	76332479483	20120	17971200	10077444472	1451520	65794754	454489400	2348920596800	2505379558400	126000
$A_1(2)$	$A_1(13)$	$E_6(3)$	$E_7(3)$	$E_8(3)$	$F_4(3)$	$G_2(4)$	${}^3D_4(3^3)$	${}^2F_4(3^2)$	${}^2B_2(2^5)$	${}^2F_4(2^3)$	${}^2G_2(3^5)$	$B_2(5)$	$C_3(7)$	$D_4(5)$	${}^2D_4(4^2)$	${}^2A_3(9)$	
20160	1092	7487687684640	14720000000000000000	147200000000000000000	5734430762856	251596800	20560831264912	32537600	244805332409	49252537	4800000	4680000	273407238	9111530800	87534471	3265920	
A_9	$A_1(17)$	$E_6(4)$	$E_7(4)$	$E_8(4)$	$F_4(4)$	$G_2(5)$	${}^3D_4(4^3)$	${}^2F_6(4^2)$	${}^2B_2(2^7)$	${}^2F_4(2^5)$	${}^2G_2(3^7)$	$B_2(7)$	$C_3(9)$	$D_5(3)$	${}^2D_4(5^2)$	${}^2A_2(64)$	
181440	2448	16167634800	100000000000000000000	1000000000000000000000	290892032340848	5409000000	47402330	34091303400	239189930264	130297600	54025791402	1389297600	34025791402	1389297600	17400203200	5315776	
A_n	$A_n(q)$	$E_6(q)$	$E_7(q)$	$E_8(q)$	$F_4(q)$	$G_2(q)$	${}^3D_4(q^3)$	${}^2F_6(q^2)$	${}^2B_2(2^{2n+1})$	${}^2F_4(2^{2n+1})$	${}^2G_2(3^{2n+1})$	$B_n(q)$	$C_n(q)$	$D_n(q)$	${}^2D_n(q^2)$	${}^2A_n(q^2)$	
$n \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	$q \geq 2$	

C_2
2
C_3
3
C_5
5
C_7
7
C_{11}
11
C_{13}
13
Z_p
C_p
p

- Alternating Groups
- Classical Chevalley Groups
- Chevalley Groups
- Classical Steinberg Groups
- Steinberg Groups
- Suzuki Groups
- Ree Groups and Tits Group*
- Sporadic Groups
- Cyclic Groups

Alternatives*
Symbol
Order†

M_{11}	M_{12}	M_{22}	M_{23}	M_{24}	$J(1), J(11)$	HJ	HJM	J_4	HS	McL	He	Ru
7920	95040	443320	10200960	244823040	175560	604800	50232960	867757104	44352000	899128000	403030720	3476264400

*For sporadic groups and families, alternate names in the upper left are other names by which they may be known. For specific non-sporadic groups these are used to indicate isomorphisms. All such isomorphisms appear on the table except the two $({}^2F_4(2) \cong {}^2F_4(2)')$.

†The size group ${}^2G_2(2)'$ is not a group of Lie type, but is the finite 2-dimensional subgroup $U_3(2)$. It is usually given binary Lie type notation.

The groups standing on the second row are the dual groups. The symbols inside give a description in the families of Suzuki groups.

*Finite simple groups are determined by their order and the following properties:
 $B_2(q)$ and $C_2(q)$ for odd $q \geq 3$
 $A_2(4)$ and $A_3(8)$ of order 240.

Sz	$O'N, O-S$	-3	-2	-1	F_4, D	Ly_5	F_5, E	$M(22)$	$M(23)$	$F_{5,2}, M(24)'$	F_8	F_4, M_1
Sz_{23}	$O'N$	C_{O_3}	C_{O_2}	C_{O_1}	HN	Ly	Th	F_{22}	F_{23}	F_{24}	B	M
448345497400	604915305920	495766456400	4230542331200	4337776806	273030	51745179	16745943	6456175145440	408947473	1235203709390	408947473	408947473

